SCHRÖDINGER'S EQUATION AND SPECIAL RADIAL ELECTRIC FIELDS

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Schrödinger's equation is exactly soluble if one considers a central potential of the type $Ar^2 + Br - D/r$ provided D takes particular well chosen values.

Many authors are interested in the search for all electric fields which allow a complete and exact solution of Schrödinger's equation. Morse and Feshbach [1] have listed convenient potentials: A/r, $A/r - B/r^2$ ($A \neq 0$), $A/r^2 + Br^2$ ($B \neq 0$). Some potentials lead to soluble equations provided the angular momentum quantum number l=0: $V=V_0\exp(-r/d)$ or $V=V_0\tanh(r/d)$. These cases excepted, Plesset [2] has shown that no exact solution in term of a finite number of elementary functions can be found if the potential is of the type $V=\sum_{k=-m}^{+n}\lambda_k r^k$ with arbitrary constant λ_k . In particular no quantization of energy exists.

Our purpose is to show for the special potential $V = Ar^2 + Br - D/r$ ($A \neq 0$) that the problem is exactly soluble if D is correctly related to A and B.

1) Spherical coordinates: $r^2 = x^2 + y^2 + z^2$. The radial part of Schrödinger's equation is:

$$\psi'' + \frac{2}{r}\psi' - \frac{l(l+1)}{r^2}\psi + \frac{2m}{\hbar^2} \left[E - Ar^2 - Br + \frac{D}{r}\right]\psi = 0,$$
(1)

$$\psi = \exp\left[-\frac{1}{2}\sqrt{\frac{2mA}{\hbar^2}}\,r^2 - \frac{1}{2}\sqrt{\frac{2m\,B^2}{A\hbar^2}}r\right]r^{-l-1}\varphi\,, \quad (2)$$

$$r\varphi'' + \left[-2\sqrt{\frac{2mA}{\hbar^2}} r^2 - \sqrt{\frac{2mB^2}{A\hbar^2}} r - 2l \right] \varphi' + \left\{ l\sqrt{\frac{2mB^2}{A\hbar^2}} + \frac{2mD}{\hbar^2} + \left[\frac{mB^2}{2A\hbar^2} + (2l-1)\sqrt{\frac{2mA}{\hbar^2}} + \frac{2mE}{\hbar^2} \right] r \right\} \varphi = 0.$$
 (3)

Eq. (3) is of the type $r\varphi'' + (ar^2 + br + c)\varphi' + (d + er)\varphi = 0$ (with c = -2l negative integer).

We have studied it elsewhere [3]: the divergent solution (at r=0, see eq. (2)) of eq. (1) is easily found as r^{-l-1} times a linear combination of Weber functions; of course we are mainly interested in the convergent quadratically integrable solution so that $\varphi=r^{2l+1}\times \text{polynomial}$ of degree n. Introducing that φ into (3) we easily deduce the two conditions to be fulfilled: e=-(n+2l+1)a and the vanishing of a determinant (more precisely a continuant [4]) of range n+1 (thus $n\geq 0$). The detailed form of that continuant is [5]:

$$\begin{split} R_k &= -2(2mA/\hbar^2)^{1/2}(k-n-2l-2) \\ S_k &= l \bigg(\frac{2mB^2}{A\hbar^2}\bigg)^{1/2} + \frac{2mD}{\hbar^2} - \bigg(\frac{2mB^2}{A\hbar^2}\bigg)^{1/2} k \end{split}$$

$$T_k = (k+1)(k-2l)$$
.

The first condition gives the energy levels which are found to be explicitly independent of D and the second condition gives the allowed D values for each value of l and n:

$$E = \sqrt{(2A \hbar^2/m)}(n+l+\frac{3}{2}) - B^2/4A$$
.

2) Cylindrical coordinates: $r^2 = x^2 + y^2$. The

same theory holds l(l+1) being simply replaced by μ^2 in eq. (1). It must be pointed out that eq. (3) is also true in this case provided one considers odd negative values of $c=(-2\mu+1)$ in place of even values as in case 1.

Conclusion. Schrödinger's equation with potential $Ar^2 + Br - D/r$ is exactly soluble only for special values of D in both spherical and cylindrical coordinates. Classically the same problem leads to elliptic integrals. Theories encountered in the general case may degenerate into elementary integrations when D is well

fitted leading to a stable periodic motion so that both classical and quantum descriptions are quite analogous.

References

- [1] P.M. Morse and H. Feshbach, Methods of theoretical physics (McGraw Hill, 1953).
- [2] M. S. Plesset, Phys. Rev. 41 (1932) 278.
- [3] A. P. Hautot, J. Math. Phys., March issue 1972.
- [4] T. Muir, A treatise on the theory of determinants (Dover S670) (1960).
- [5] A. P. Hautot, Bull. Soc. Roy. Sci. Liège 38 (1969)

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